

*Linear stability and Morse index
for the figure-eight and $k=5$ slalom
solutions
under homogeneous potential*

Toshiaki Fujiwara, Hiroshi Fukuda, and Hiroshi Ozaki

*Celestial Mechanics and N-Body Problem
2018/07/05 AIMS Taipei*

Three-body choreographies and Continuations

$$L = \frac{1}{2} \sum_k \left| \frac{dq_k}{dt} \right|^2 + \frac{1}{\alpha} \sum_{i,j} \frac{1}{|q_i - q_j|^\alpha}$$

$\alpha = 1$: Newton potential

$$q_0(t) = q(t), q_1(t) = q(t + T/3), q_2(t) = q(t + 2T/3)$$

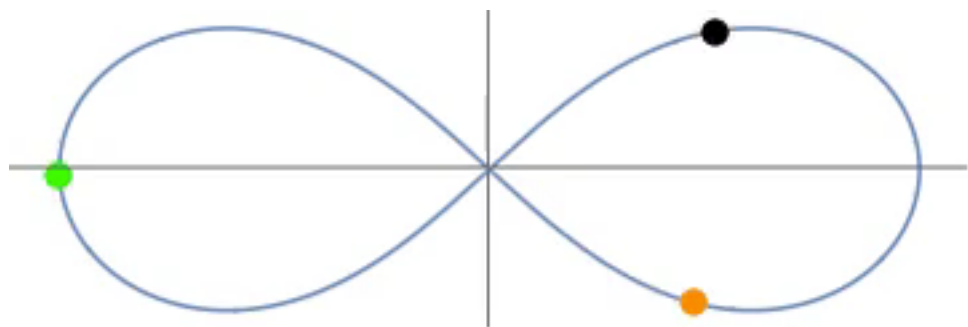


figure-eight solution

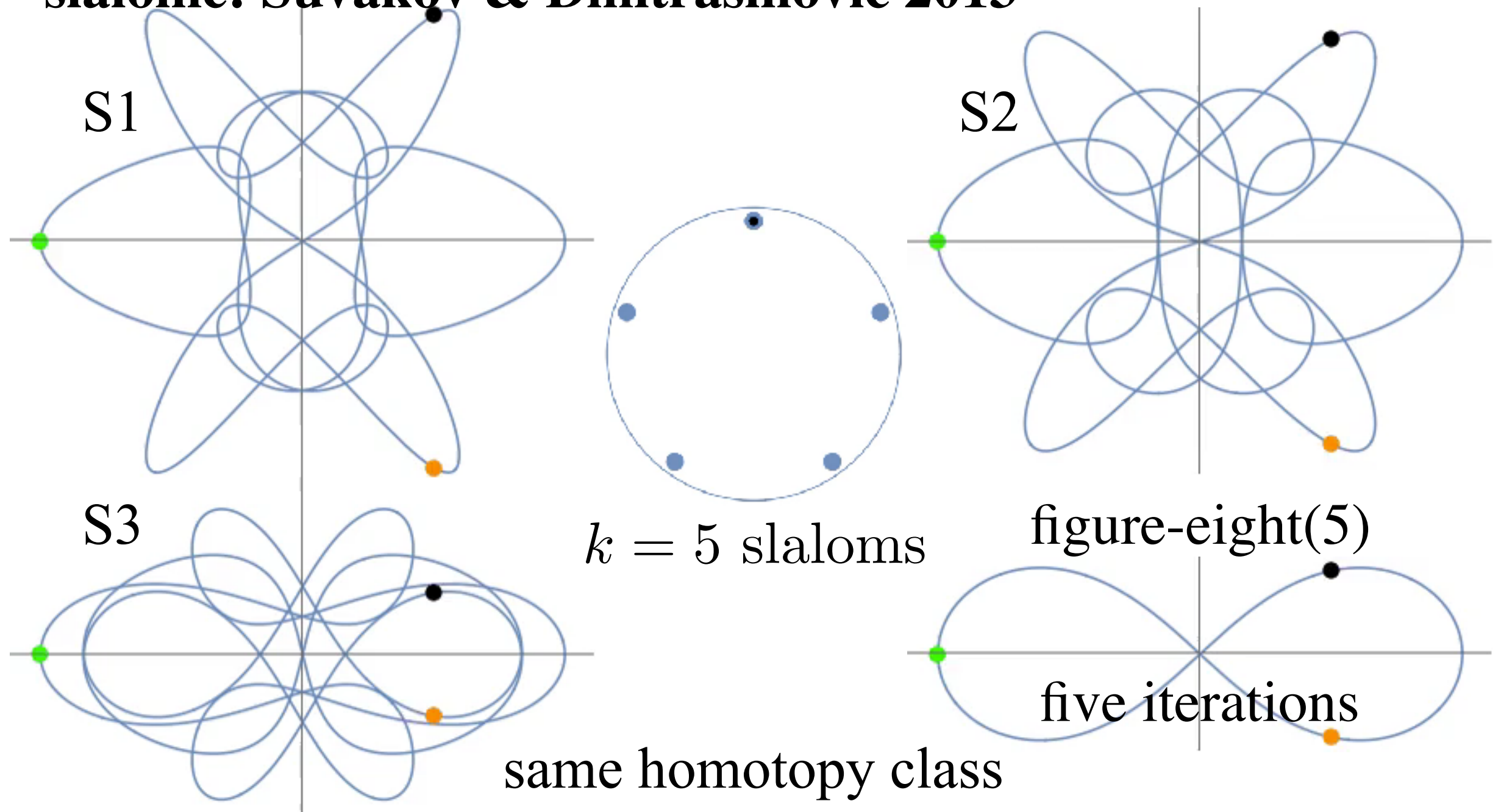
C. Moore 1993,

A. Chenciner and R. Montgomery 2000

Figure-eight and slalom solutions

figure-eight: Moore 1993, Chenciner & Montgomery 2000

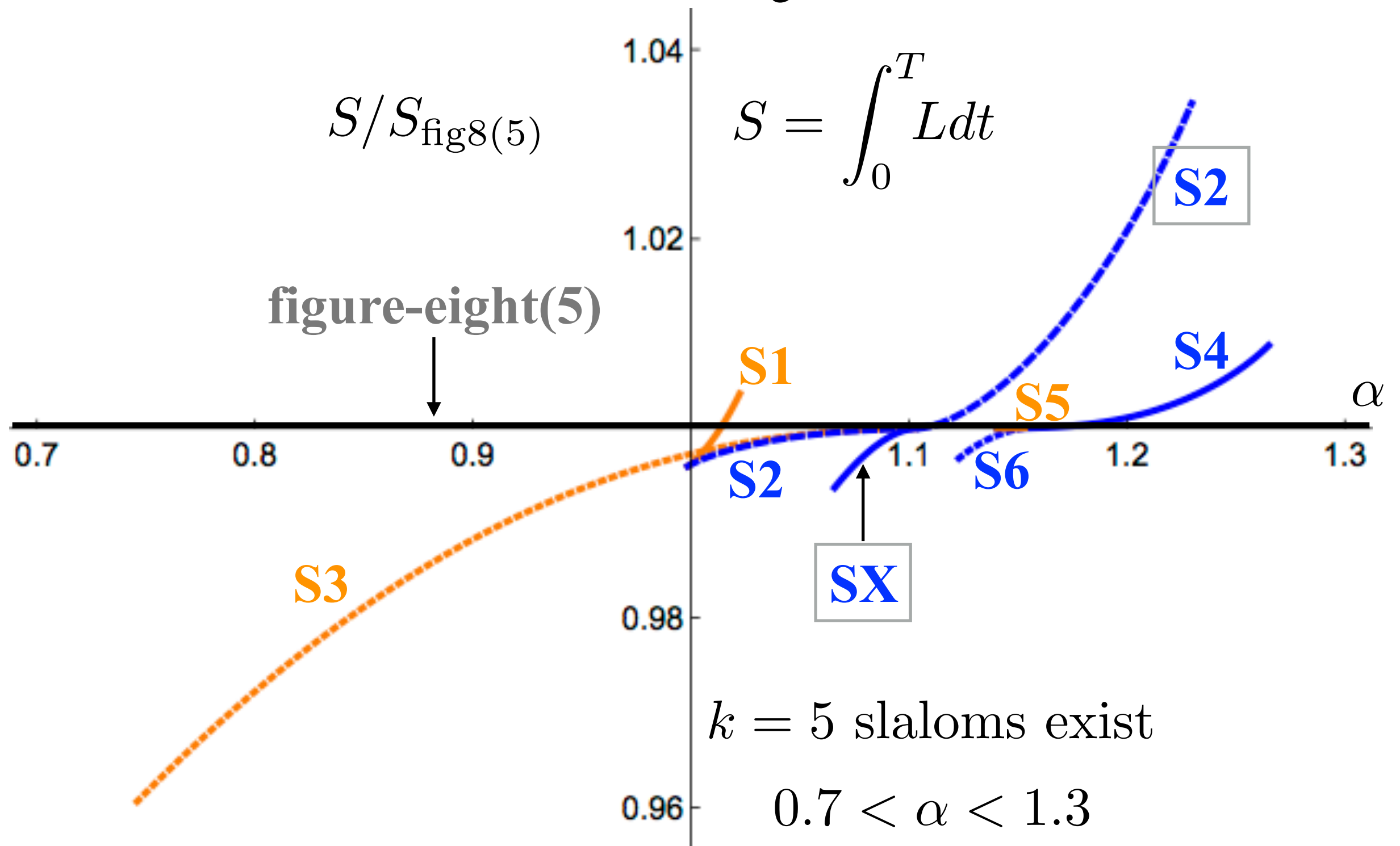
slalome: Šuvakov & Dmitrašinović 2013



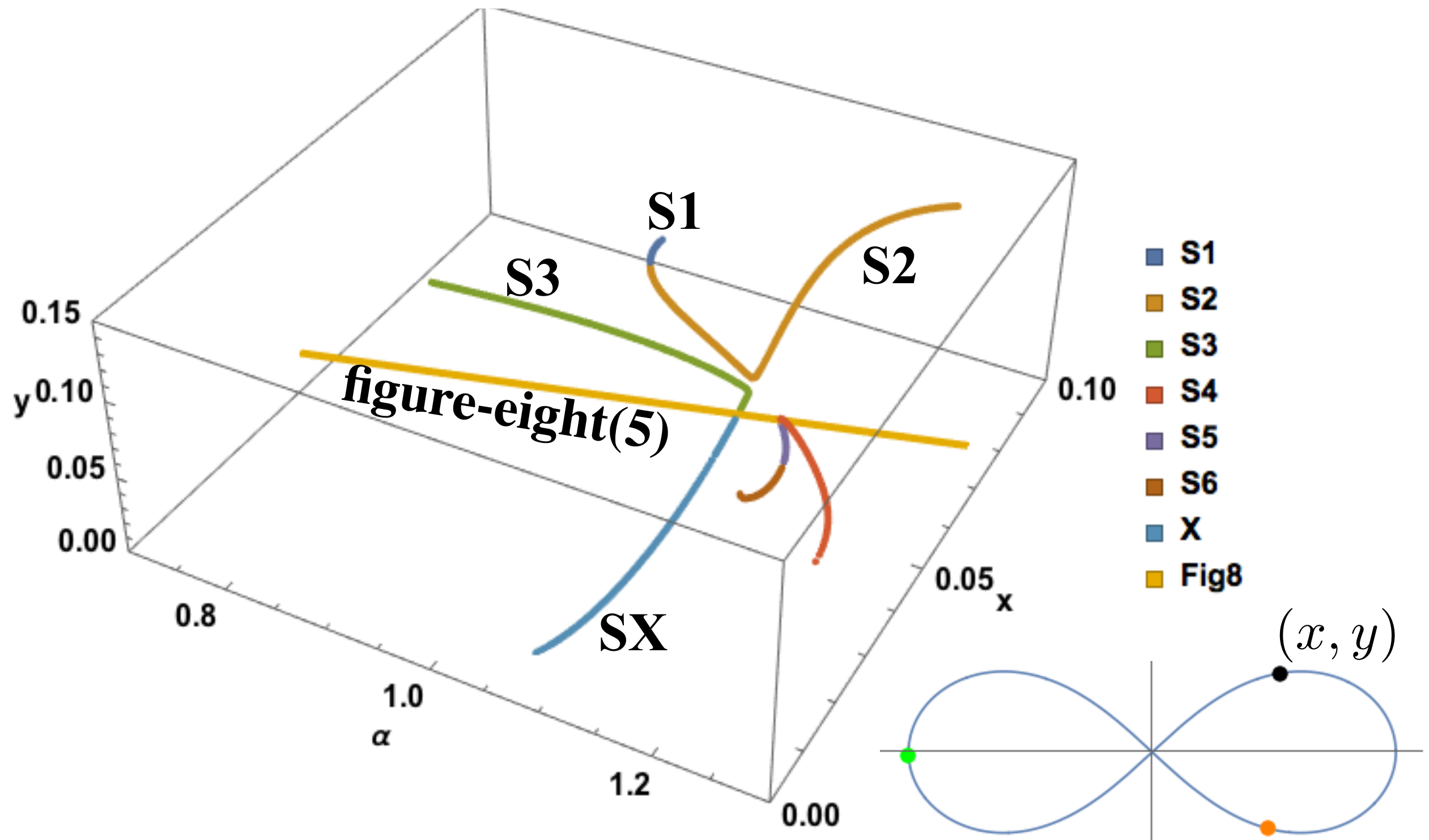
Slalom Solutions

- Šuvakov M.
 - Numerical search for periodic solutions in the vicinity of the figure-eight orbit: slaloming around singularities on the shape sphere, *Celest. Mech. Dyn. Astron.* 119, 369–377 (2014)
- Šuvakov M., Dmitrašinović V.
 - Three classes of Newtonian three-body planar periodic orbits, *Phys. Rev. Lett.* 110(11), 114301 (2013)
- Šuvakov M., Dmitrašinović V.
 - A guide to hunting periodic three-body orbits, *Am. J. Phys.* 82, 609–619 (2014)
- Šuvakov M., Shibayama M.
 - Three topologically nontrivial choreographic motions of three bodies, *Celest. Mech. Dyn. Astron.* 124, 155–162 (2016)

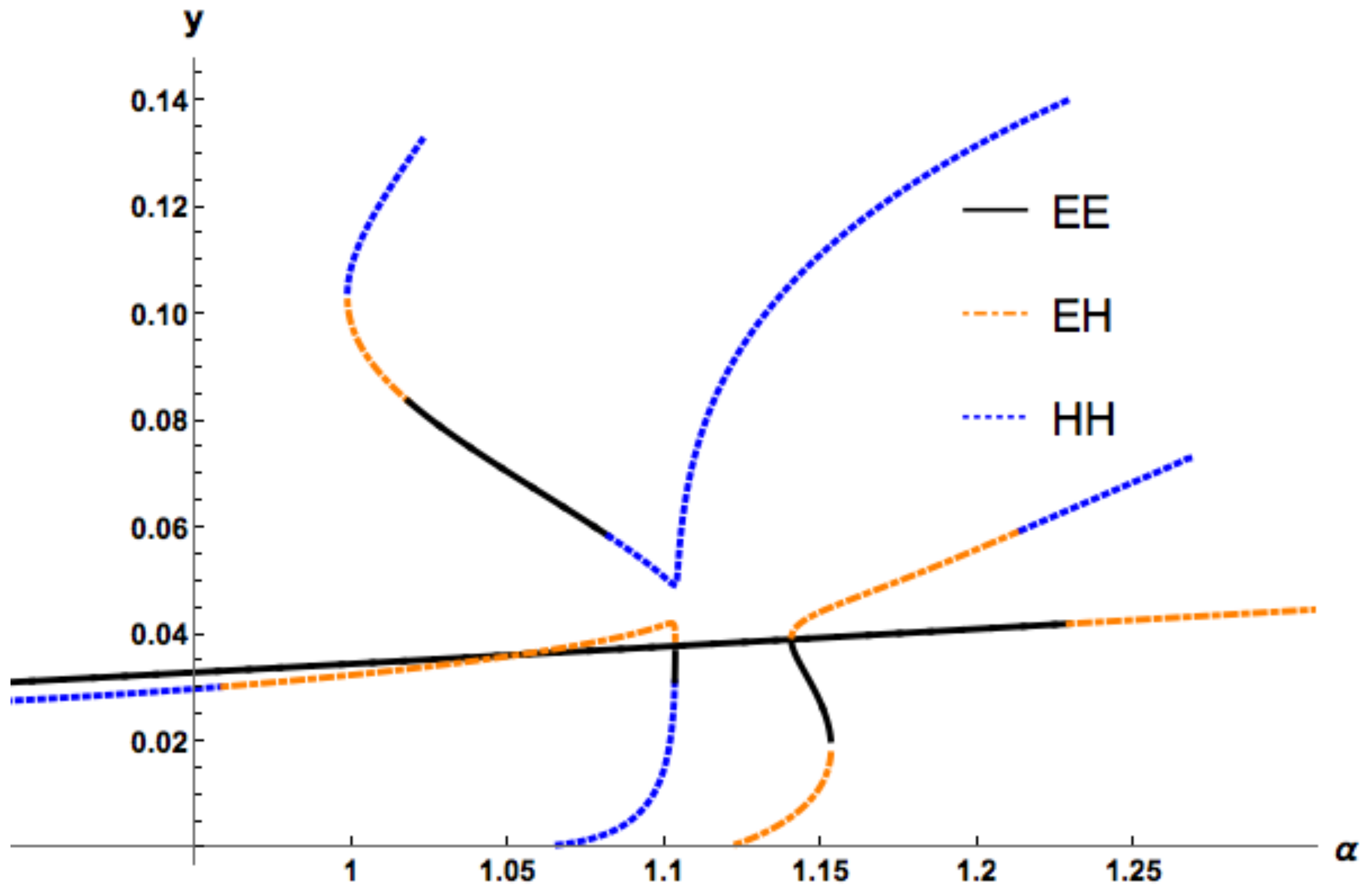
Continuation of solutions



continuation of solutions



Linear stability



Continuation of
known slalom solutions $k=5$,
S1, S2, S3



***Period 5 bifurcations
from the figure-eight***

Linear stability
Second derivative of action for figure-eight(5)



other bifurcated choreographic solutions?

Yes!

Linear stability, Floquet matrix

small variation $\delta q(t), \delta p(t) \in \mathbb{R}^6$
around a periodic solution $q(t), p(t)$
with $q(t+T) = q(t), p(t+T) = p(t)$

equation of motion for $\Psi(t) = \begin{pmatrix} \delta q(t) \\ \delta p(t) \end{pmatrix}$, $\frac{d}{dt}\Psi(t) = B(t)\Psi(t)$

Green function $G(t)$, $\frac{d}{dt}G(t) = B(t)G(t) \Rightarrow \Psi(t) = G(t)\Psi(0)$

Floquet matrix $M = G(T)$

$$M\Psi = \mu\Psi$$

μ : characteristic multiplier

A curved arrow originates from the expression $\frac{\partial^2 V}{\partial q_i \partial q_j}$ and points to the matrix $B(t)$ in the equation of motion $\frac{d}{dt}\Psi(t) = B(t)\Psi(t)$. The arrow is labeled with δ_{ij} at its tip and has a '0' at its tail.

characteristic multipliers μ and linear stability

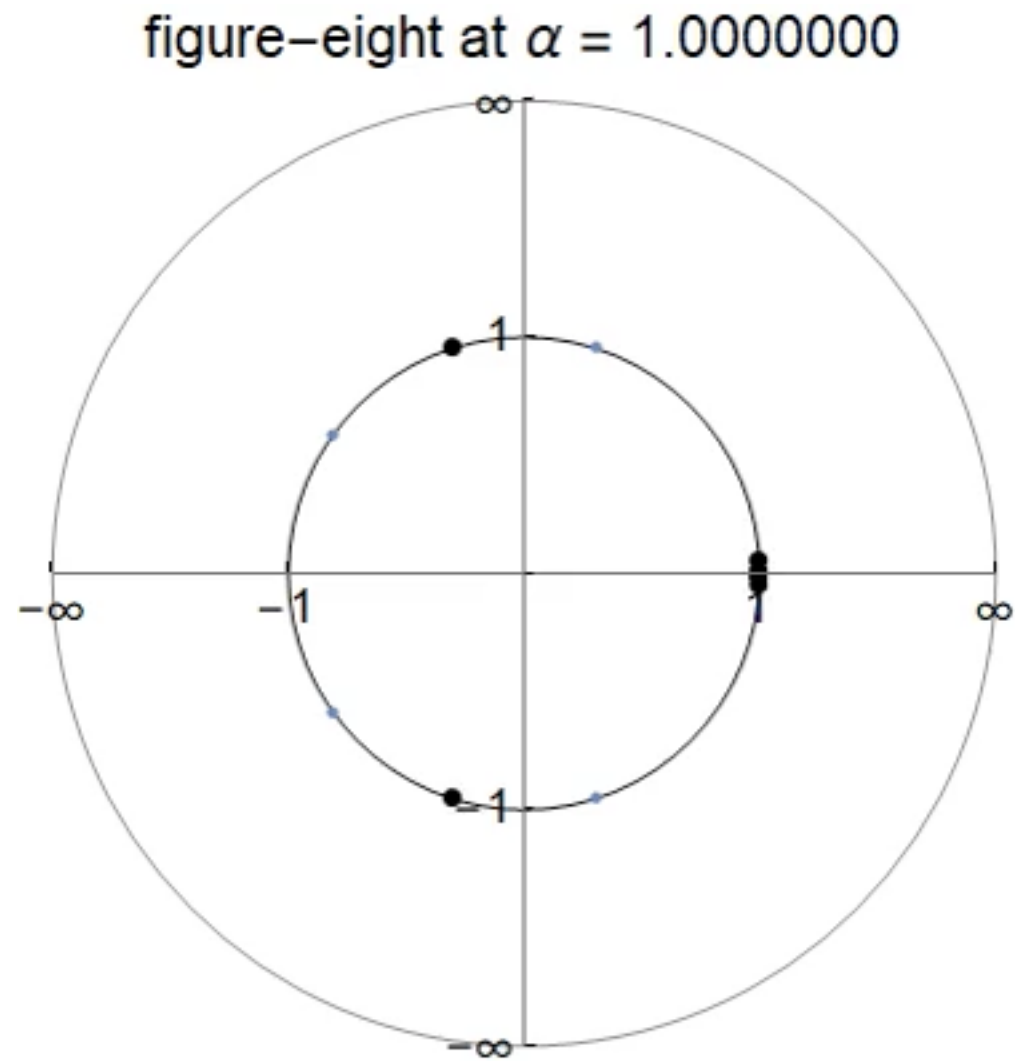
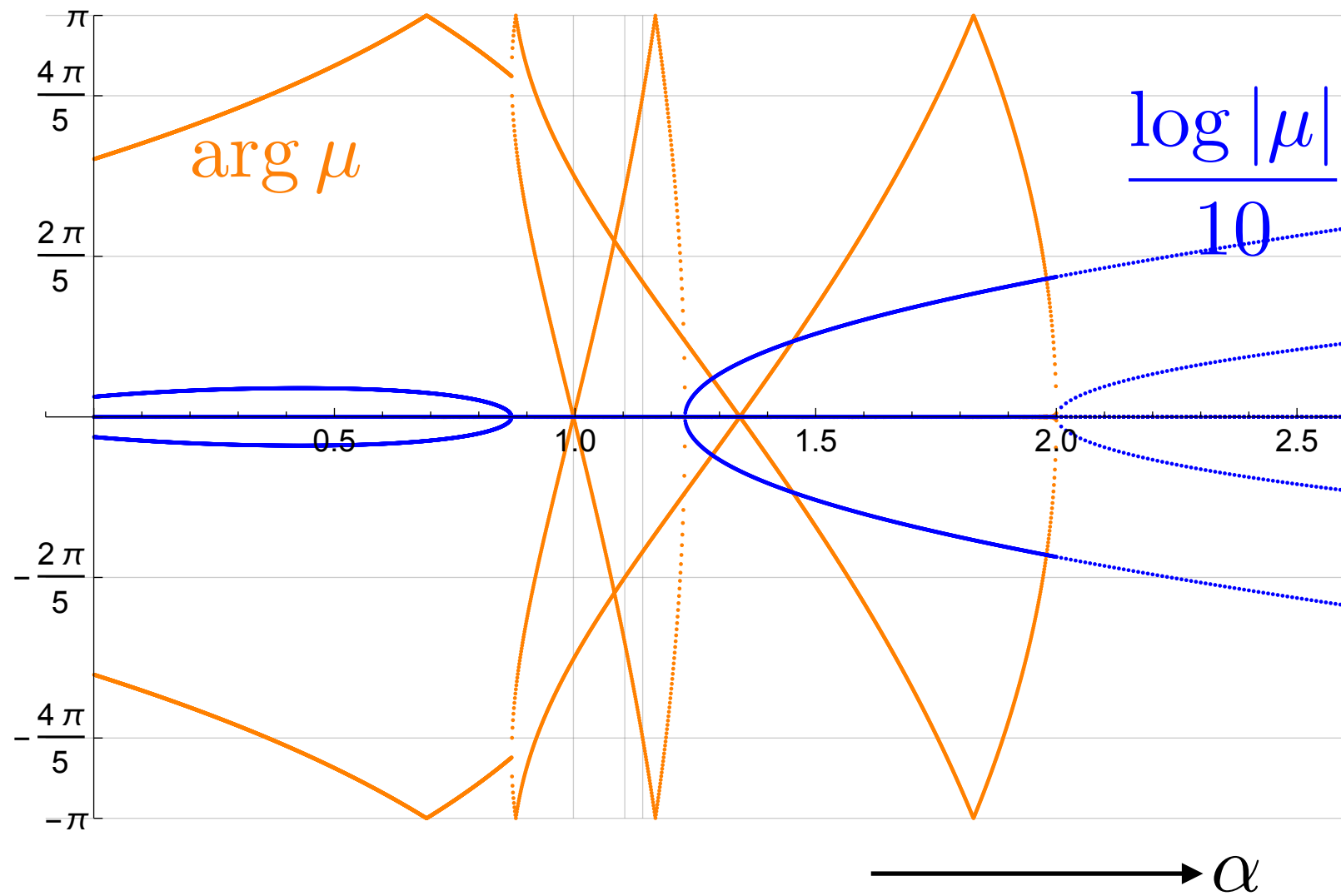
$$M\Psi = \mu\Psi$$

8 trivial $\mu = 1$, 4 non-trivial μ

M : symplectic $\Rightarrow \mu, 1/\mu, \mu^*, 1/\mu^*$

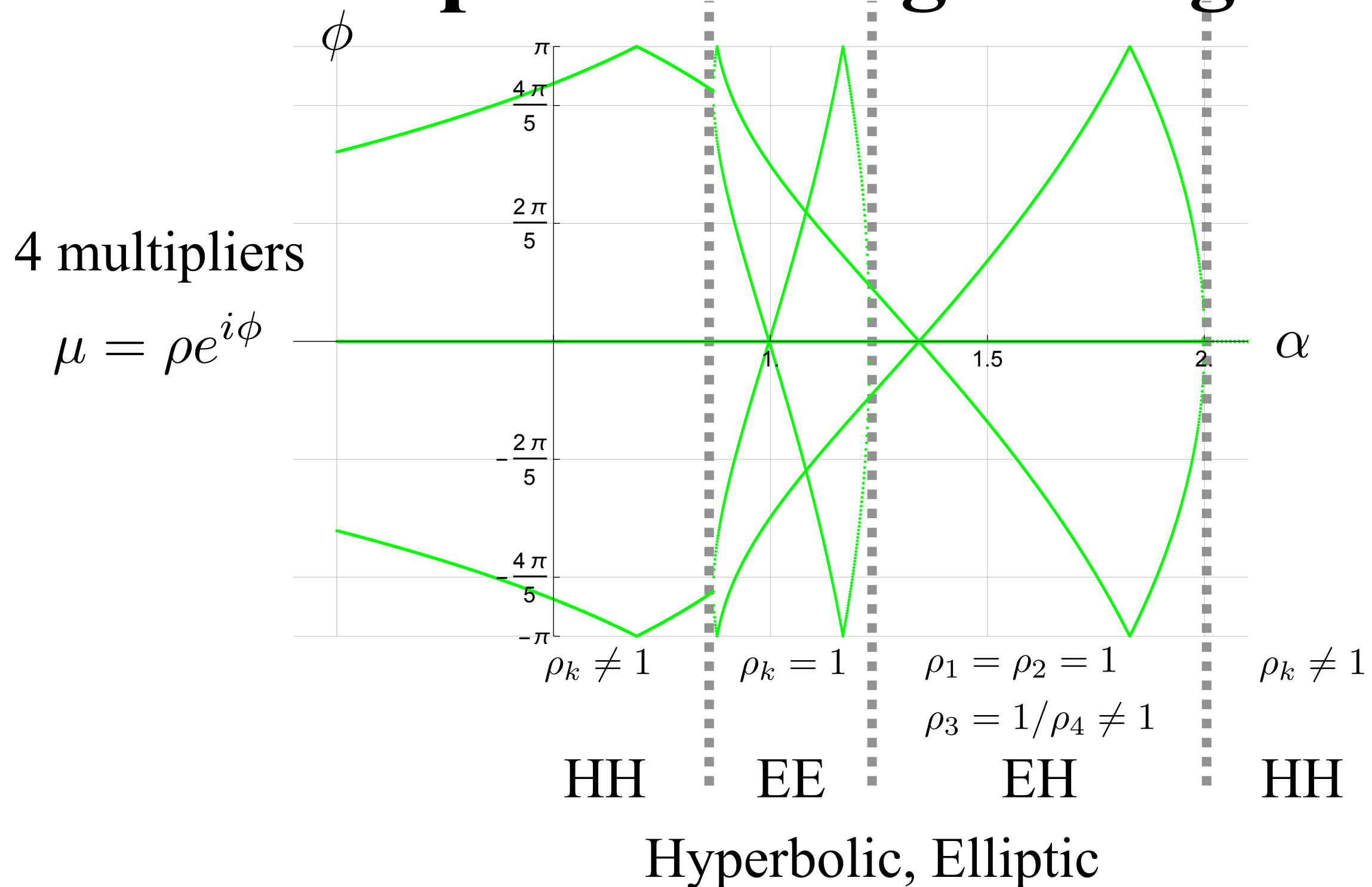
$$\begin{cases} |\mu| = 1 \rightarrow \text{E: elliptic,} \\ |\mu| \neq 1 \rightarrow \text{H: hyperbolic} \end{cases}$$

characteristic multipliers for figure-eight(1)



$$\mu_i \in \mathbb{R}, i = 1, 2, 3, 4$$

arguments of characteristic multipliers for figure-eight



Hessian of action

Equal mass planar three-body problem

$$L = \frac{1}{2} \sum_{\ell} \left(\frac{dq_{\ell}}{dt} \right)^2 + U, \quad U = \sum_{i \neq j} V(|q_i - q_j|)$$

$$S[q + \delta q] = S[q] + \underbrace{\delta S[q]}_{=0} + \frac{1}{2} \int_0^T dt \sum_{i,j} \delta q_i \left(\underbrace{-\delta_{ij} \frac{d^2}{dt^2} + \frac{\partial^2 U}{\partial q_i \partial q_j}}_{= \mathcal{H}: \text{Hessian}} \right) \delta q_j$$

$$\mathcal{H}\Psi = \lambda\Psi, \quad \Psi = \begin{pmatrix} \delta q_0 \\ \delta q_1 \\ \delta q_2 \end{pmatrix}, \quad \delta q_{\ell} = \begin{pmatrix} \delta q_{\ell x} \\ \delta q_{\ell y} \end{pmatrix} \in \mathbb{R}^2.$$

$\delta q_{\ell}(t + T) = \delta q_{\ell}(t)$: only impose periodicity for variations

$$T = 5T_{\text{figure-eight}}$$

symmetries

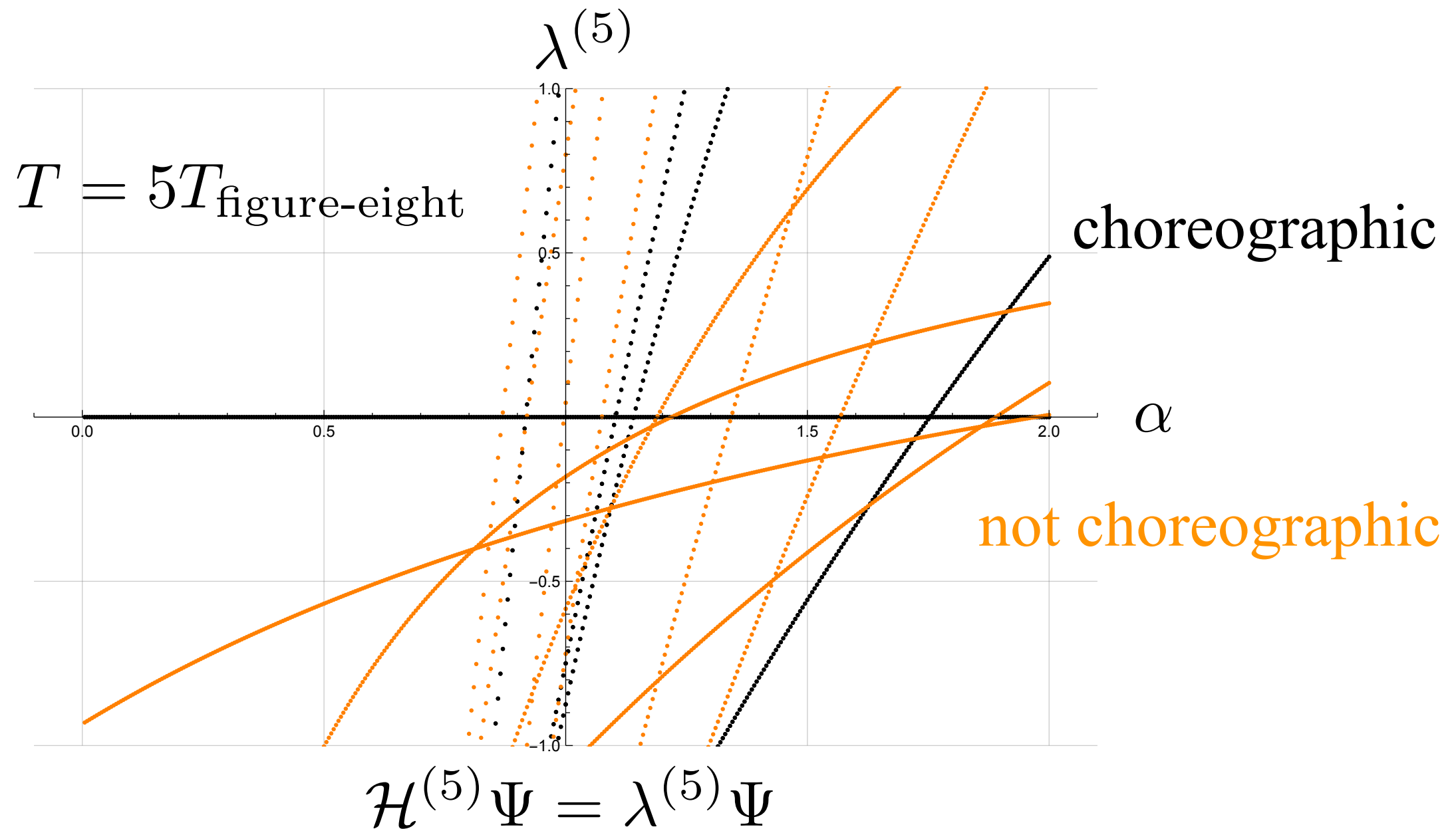
Since the figure-eight solution has symmetries

$$\begin{cases} \mathcal{C} & : \text{choreographic} \\ \mathcal{M} & : t \rightarrow t + T/2 \\ \mathcal{T} & : t \rightarrow -t \end{cases}$$

So, the Hessian around the figure-eight is
invariant under the symmetries.

Then, the eigenfunction $\mathcal{H}^{(5)}\Psi = \lambda\Psi$
can be classified by $\{\mathcal{C}, \bar{\mathcal{C}}\}, \mathcal{M}^\pm, \mathcal{T}^\pm$

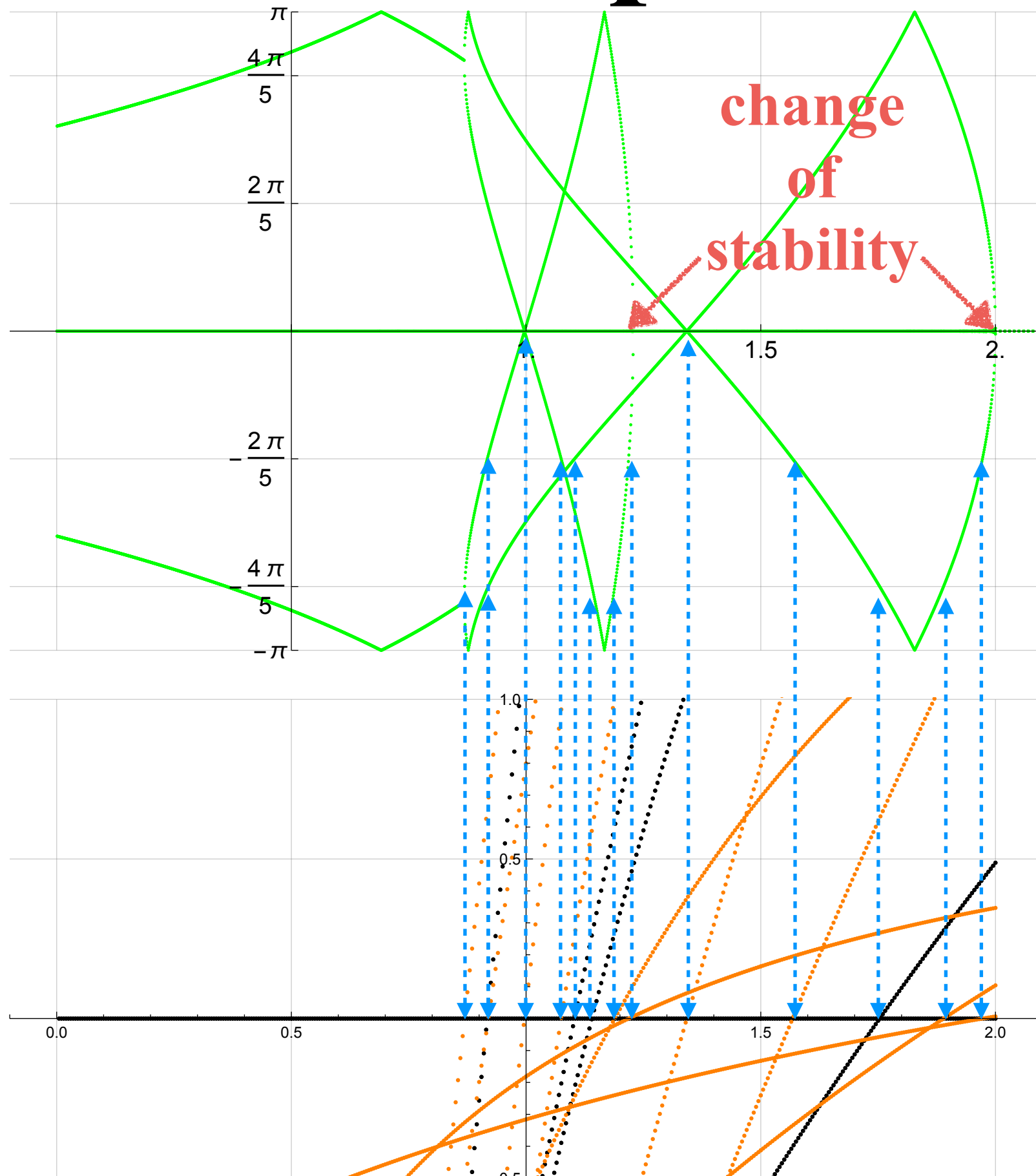
eigenvalues of Hessian of Action



every eigenvalues here are **doubly degenerated**

1 to 1 correspondence

$\phi = \pm \frac{2n\pi}{5}$
 1 to 1
 with
 2 exceptions
 $\lambda^{(5)} = 0$



1 to 1

characteristic multipliers: $\mu^5 = 1$

bifurcation + **change of stability**

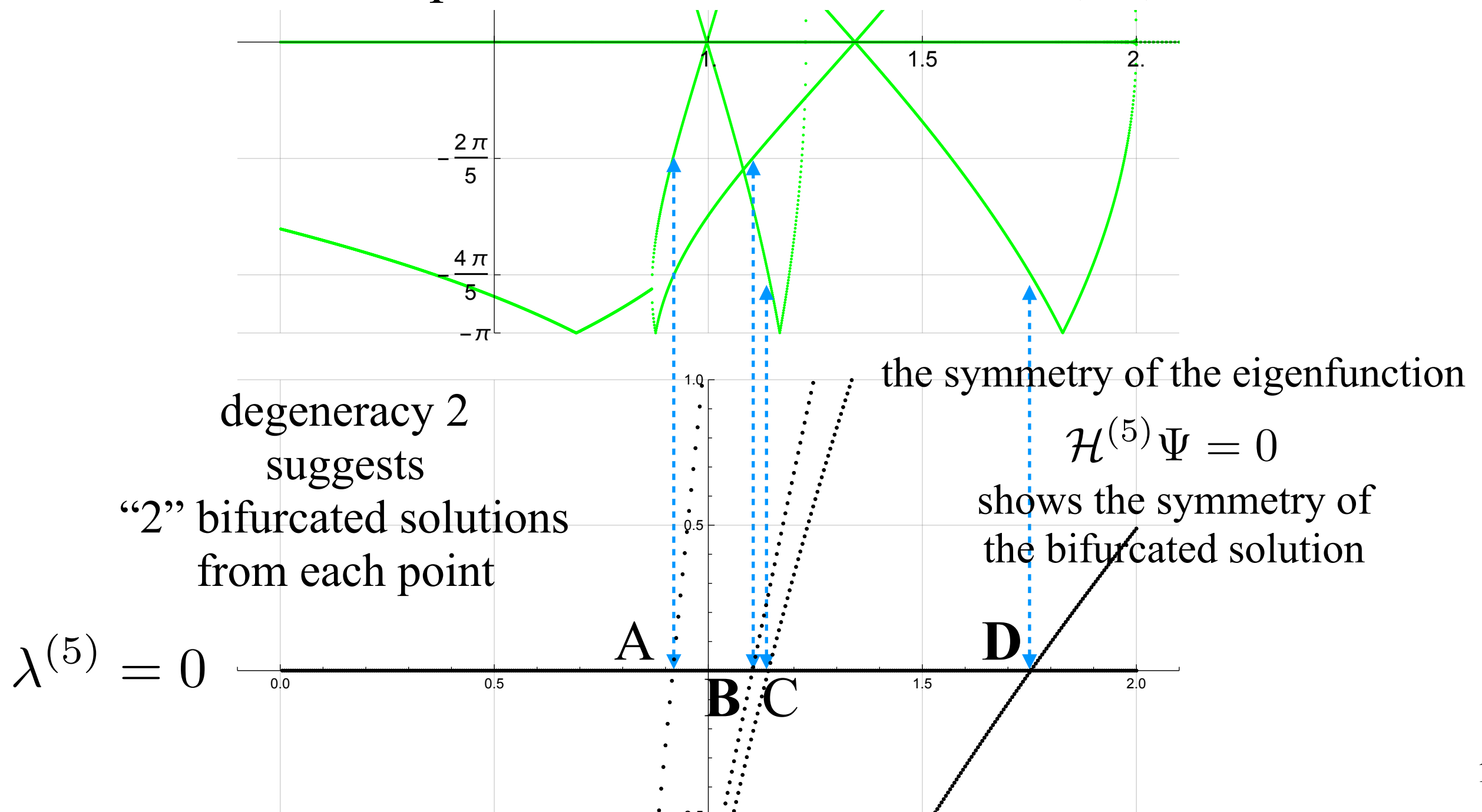
eigenvalues of Hessian of Action: $\lambda^{(5)} = 0$
multiplicity of $\lambda^{(5)} = 0$ = “number” of bifurcated solutions

This correspondence is confirmed for figure-eight(1)
under Lennard-Jones and homogeneous potential

Question: for figure-eight(5) ?

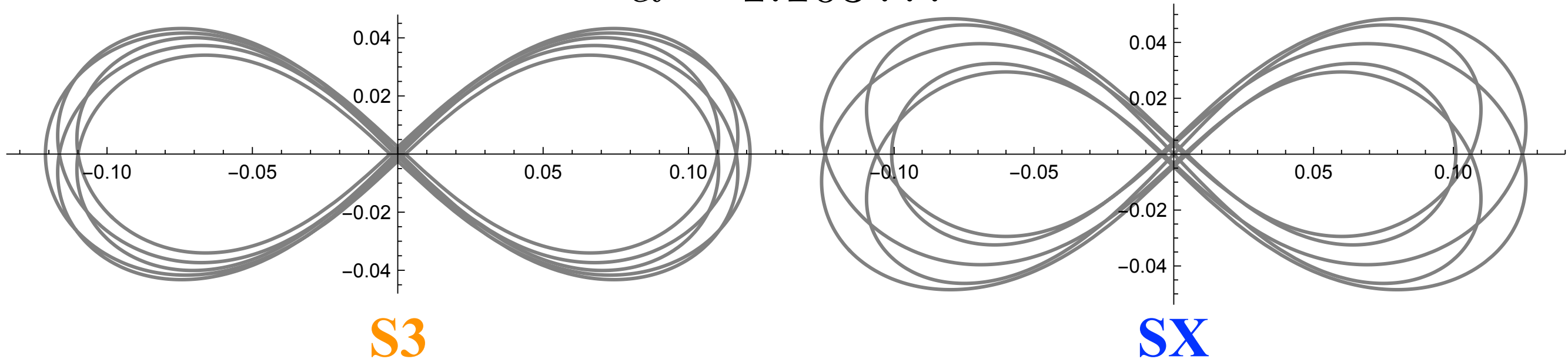
4 $\lambda=0$ points for choreographic variation

The correspondence is confirmed for B, C, D.



bifurcated solutions near the point B

$$\alpha \sim 1.103 \dots$$



they are choreographic and
keep the same symmetry as the figure-eight

Eigenvalues of H near branch point B

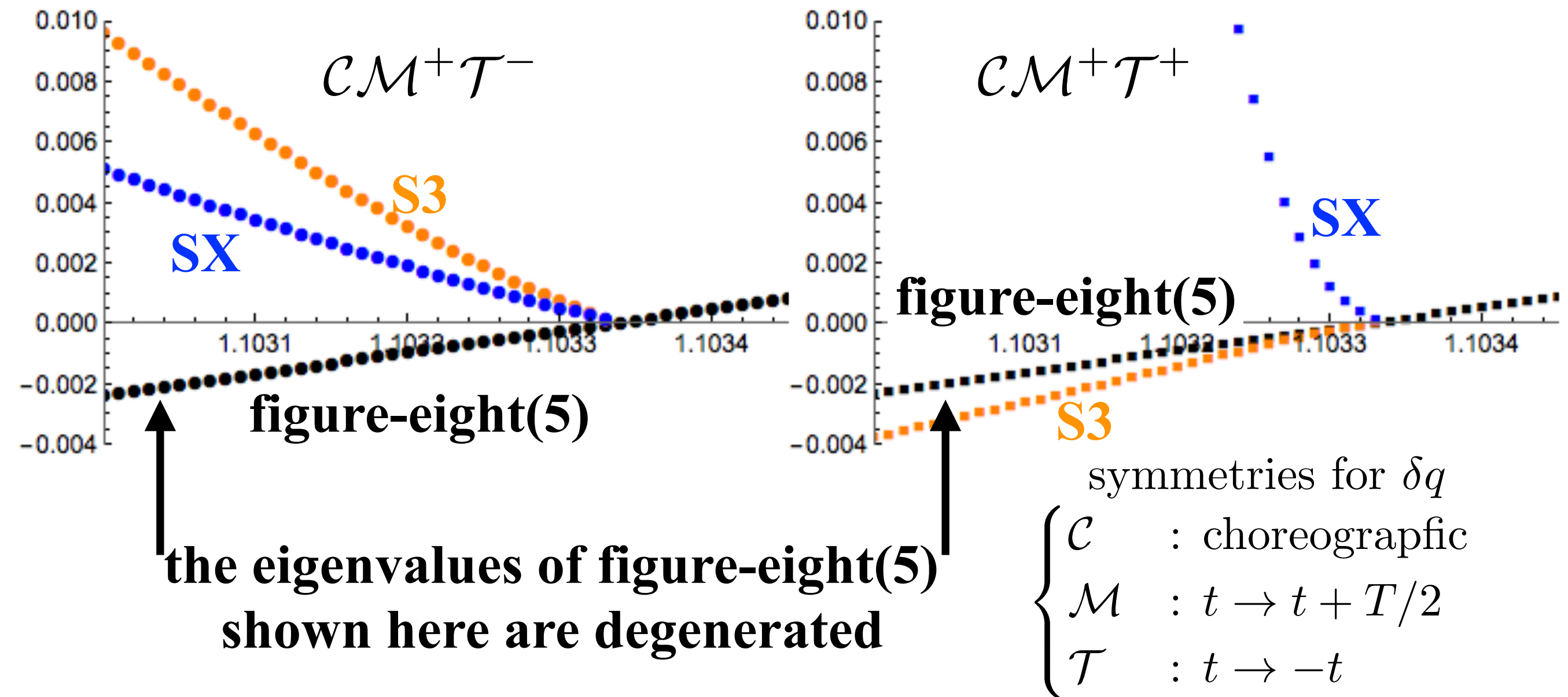
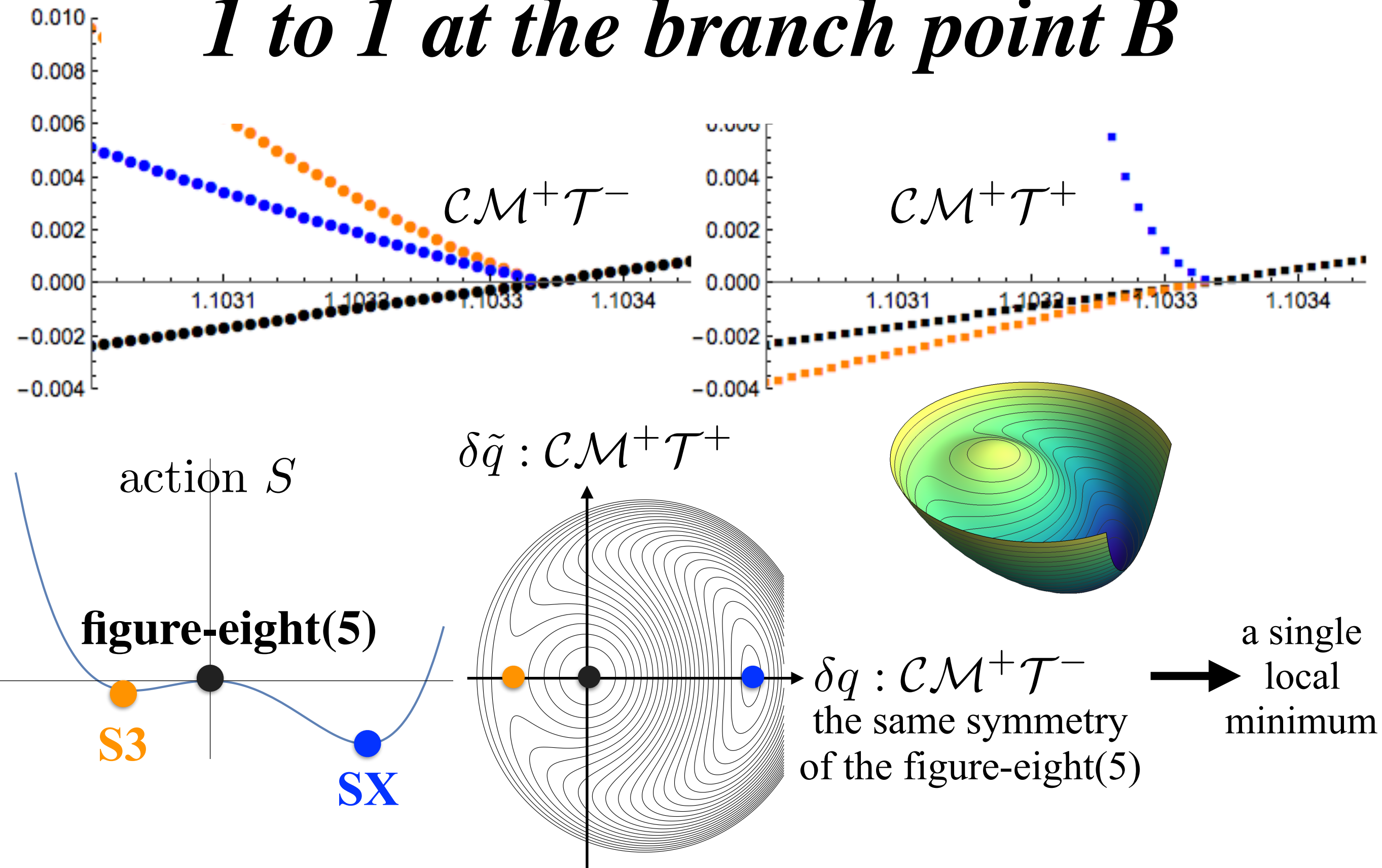
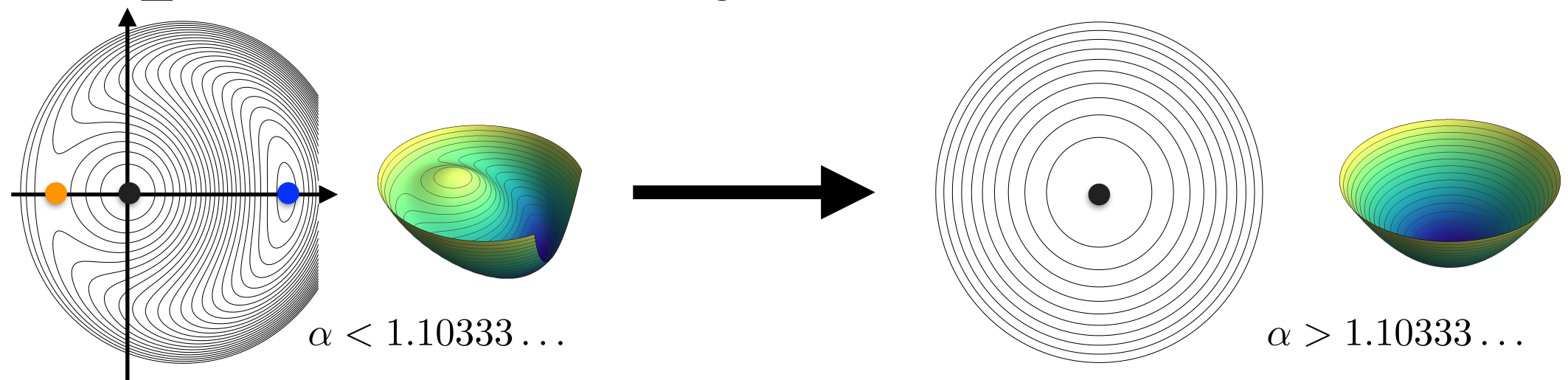


figure-eight q has $\mathcal{CM}^+ \mathcal{T}^-$

1 to 1 at the branch point B



Speculation for the action



Images are just a speculation for the action based on the behaviour of eigenvalues of Hessian of the action shown in the previous slide. We don't calculate the action directly. The XY directions represent the two eigenfunction of the Hessian for figure-eight(5),

$$\mathcal{H}^{(5)} e_k = \lambda^{(5)} e_k, k = 1, 2.$$

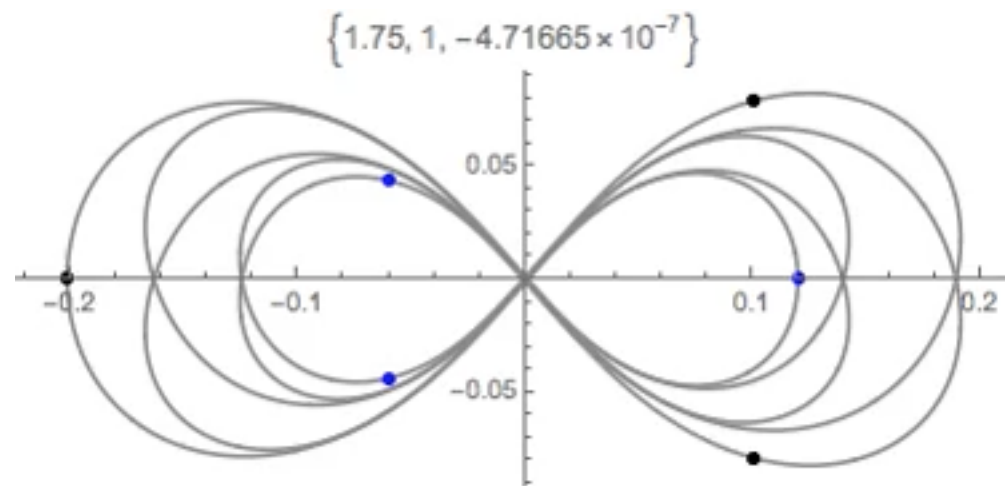
Since the eigenvalues for the figure-eight that we are noticing is doubly degenerated, we have two eigenfunctions e_1 and e_2 that are represented by δq and $\delta \tilde{q}$ in the previous slide.

For $\alpha < 1.10333\dots$, the eigenvalues for figure-eight(5) are (both) negative, the eigenvalues for S_3 are positive and negative for each direction (see the previous slide) and the eigenvalues for S_X are both positive. Therefore, the point figure-eight is a local maximum, S_3 is a saddle, and point S_X is a local minimum.

On the other hand, for $\alpha > 1.10333\dots$, the solutions S_3 and S_X vanish, and only the figure-eight(5) has (doubly degenerated) positive eigenvalues. Therefore, the action has simple local minimum at the figure-eight.

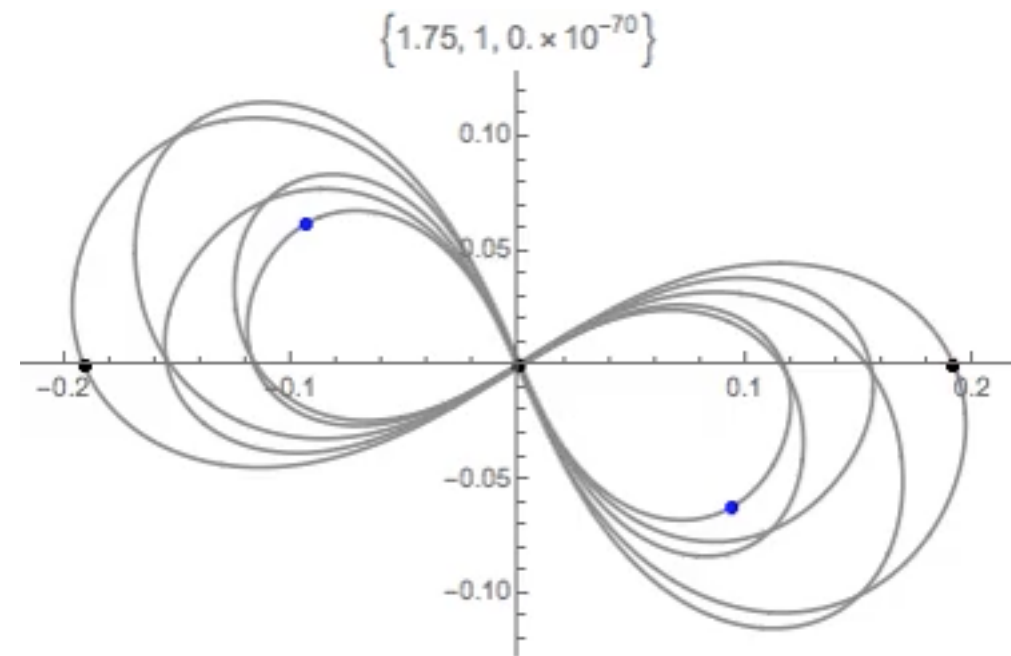
near the point D

$$\alpha \sim 1.75 \dots$$



$$c \neq 0$$

X-axis symmetric,
not Y-axis

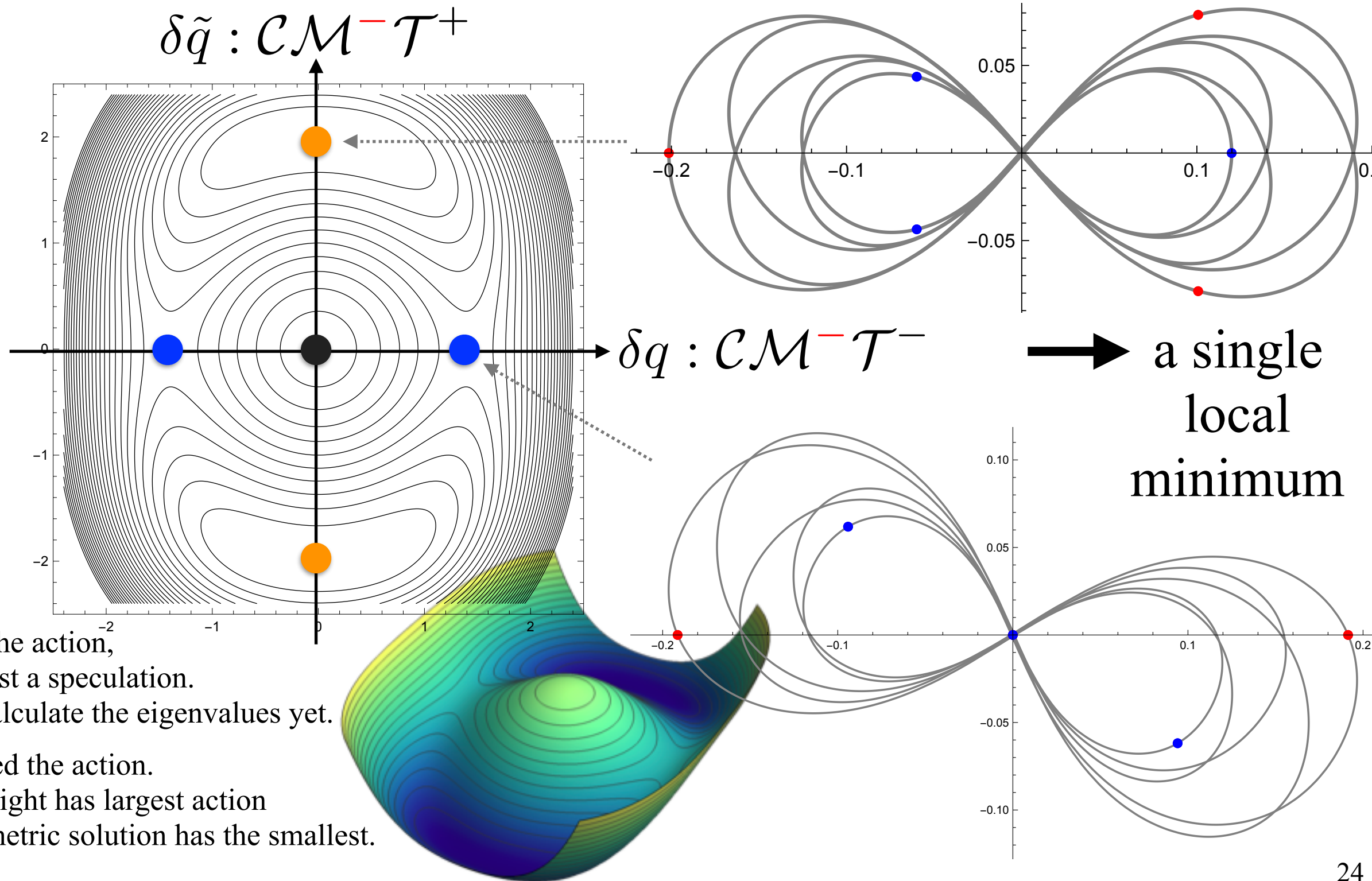


$$c = 0$$

Origin symmetric

choreographic
but the figure-eight symmetry is broken

1 to 1 at the branch point D



summary

not “proved”

We **observed** the following properties for the bifurcations from figure-eight(5);

- characteristic multiplier $\mu=1 \rightarrow$ bifurcation or change of stability
- eigenvalues of Hessian $\lambda=0$, $\mathcal{H}^{(5)}\Psi = 0 \Leftrightarrow$ bifurcation (1 to 1)
- multiplicity of $\lambda=0 \Leftrightarrow$ “number” of bifurcated solutions
- symmetry of the eigenfunction δq predict the symmetry of bifurcated solution $q+\delta q$
- eigenvalues of Hessian for the figure-eight(5) and the bifurcated solutions suggest the behaviour of action in function space

Summary

