Synchronised Similar Triangles for Three Body Orbit with Zero Angular Momentum

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> June 1 SPT 2004 Cala Gonone, Sardinia, Italy

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http://arxiv.org/abs/**math-ph/0404056** http://**www.clas.kitasato-u.ac.jp/~fujiwara/**nBody/nbody.html

Three-body figure-eight choreography

- C. Moore (1993): found numerically
- A. Chenciner and
 R. Montgomery (2000):
 proved the existence



- C. Simó (2000):
 found a lot of N-body choreographies numerically
- Barutello, Ferrario, Terracini, Chen, Shibayama, ...

Three-body figure-eight choreography

Treasure to be a state of the state of

$$i = 1, 2, 3, \ m_i = 1$$
$$\ddot{q}_i = \sum_{j \neq i} \frac{q_j - q_i}{|q_j - q_i|^3},$$
$$\begin{cases} q_1(t) &= q(t), \\ q_2(t) &= q(t + T/3), \\ q_3(t) &= q(t + 2T/3) \end{cases}$$

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$$\sum_{i} q_i = 0, \ L = \sum_{i} q_i \wedge \dot{q}_i = 0.$$

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Figure-eight has zero angular momentum

Why L = 0?

Total angular momentum is conserved. Therefore,



Then, what does L = 0 mean?

Three tangents theorem (FFO 2003) L=0

Theorem (Three Tangents). If $\sum_i p_i = 0$ and $\sum_i q_i \wedge p_i = 0$, then three tangents meet at a point.



Three tangents theorem

Theorem (Three Tangents). If $\sum_i p_i = 0$ and $\sum_i q_i \wedge p_i = 0$, then three tangents meet at a point.

Proof. Let C_t be the crossing point of two tangent lines p_1 and p_2 .

$$\therefore (q_1 - C_t) \land p_1 = 0, (q_2 - C_t) \land p_2 = 0$$

On the other hand, $\sum_i p_i = 0$, $\sum_i q_i \wedge p_i = 0 \Rightarrow \sum_i (q_i - C_t) \wedge p_i = 0$.

$$\therefore (q_3 - C_t) \land p_3 = 0. \qquad \Box$$

 C_t : the "Center of Tangents"

Centre of force for three body orbit dL/dt=0

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Since

$$\sum_{i} f_i = 0, \quad \frac{dL}{dt} = \sum_{i} q_i \wedge f_i = 0,$$

force vector from each bodies meet at a point C_f : the Centre of force. (Schiaparelli, Wintner)

Three tangents theorem L=0

- Shape of the orbit of Figure Eight x(t) and the orbit C(t) are still unknown.
- Three Tangents Theorem gives a criterion for the orbit.
- For example ...



The simplest curve: 4-th order polynomial

$$x^4 + \alpha x^2 y^2 + \beta y^4 = x^2 - y^2$$



Candidate: Lemniscate and its scale transform

$$\begin{aligned} (x^2+y^2)^2 &= x^2-y^2 \\ x &\to \mu x, y \to \nu y \end{aligned}$$

Three-body choreography on the lemniscate (FFO 2003)

Choreograpgy on the Lemniscate

$$q(t) = \left(\frac{\operatorname{sn}(t)}{1 + \operatorname{cn}^2(t)}, \frac{\operatorname{sn}(t)\operatorname{cn}(t)}{1 + \operatorname{cn}^2(t)}\right) \text{ with } k^2 = \frac{2 + \sqrt{3}}{4}$$

$$\begin{cases} q_1(t) = q(t), \\ q_2(t) = q(t + T/3), \\ q_3(t) = q(t + 2T/3) \end{cases}$$

satisfies the equation of motion $\ddot{q}_i = -\frac{\partial}{\partial q_i}U$ with

$$U = \sum_{i < j} \left(\frac{1}{2} \ln r_{ij} - \frac{\sqrt{3}}{24} r_{ij}^2 \right).$$

Figure-eight solution under homogeneous potential

$$V_{\alpha} = \begin{cases} \alpha^{-1} r^{\alpha} & \text{for } \alpha \neq 0\\ \log r & \text{for } \alpha = 0 \end{cases}$$

Numerical evidence Moore: Exist for $\alpha < 2$ CGMS: Exist for $\alpha < 0$ and Stable $\alpha = -1 \pm \epsilon$



Figure-eight for $\alpha = -2$ has $I = \sum_{i} q_i^2 = \text{const.}$

Evolution of moment of inertia

$$\sum_{i} m_{i}q_{i} = 0, \quad M = \sum_{i} m_{i}, \quad I = \sum_{i} m_{i}q_{i}^{2},$$

$$K = \sum_{i} \dot{q}_{i}^{2}, \quad V_{\alpha} = \frac{1}{\alpha} \sum_{i < j} m_{i}m_{j}r_{ij}^{\alpha}, \quad H = \frac{1}{2}K + V_{\alpha}.$$

$$\Rightarrow \frac{d^{2}I}{dt^{2}} = 2K - 2\alpha V_{\alpha} = 4E - 2(2 + \alpha)V_{\alpha}.$$
Lagrange-Jacobi identity
For $\alpha = -2: \quad \frac{d^{2}I}{dt^{2}} = 4E \Rightarrow I = 2Et^{2} + c_{1}t + c_{2}.$

$$\therefore \text{ If } I \neq 0 \text{ and } I \neq \infty \text{ then } E = 0, \quad c_{1} = 0, \quad I = \text{ const.}$$
Figure-eight under $1/r^{2}$ has $I = \text{constant.}$

What does this mean?

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Three normals theorem for dI/dt=0 orbit



This theorem holds for general masses m_i .

 C_n : the "Center of Normals"

Circumcircle theorem for L=0, dI/dt=0 orbit

Theorem (CircumCircle). If $\sum_{i} p_{i} = 0, \sum_{i} q_{i} \wedge p_{i} = 0$ and $\sum_{i} q_{i} \cdot p_{i} = 0$, then C_{t} and C_{n} are the end points of a diameter of the circumcircle for the triangle $q_{1}q_{2}q_{3}$.

Proof. Angles $C_t q_i C_n$ are 90 degrees for i = 1, 2, 3.

This theorem holds for any masses m_i .



Centres for figure-eight solution under 1/r^2 potential



Synchronised similar triangles for figure-eight under 1/r^2



Two triangles are inversely congruent. Because ...

Similar triangles for L=0, dI/dt=0 orbit.

Theorem (Similar Triangles). If $\sum_i p_i = 0$, $\sum_i q_i \wedge p_i = 0$ and $\sum_i q_i \cdot p_i = 0$, then triangle whose vetices are q_i and triangle whose perimeters are p_i are similar with reverse orientation.

Proof.Look at theangles yellow coloredand red colored.It is obvious.

Remark: This theorem holds for any masses m_i



Ratio of magnification (L=0, dI/dt=0)

$$\kappa(t) = \frac{|p_k|}{|q_i - q_j|} = \sqrt{\frac{m_1 m_2 m_3 K}{MI}}$$

$$\therefore \frac{\kappa(t)^2}{m_1 m_2 m_3} = \frac{p_k^2 / m_k}{m_i m_j (q_i - q_j)^2} = \frac{K}{MI}$$
where $K = \sum_i \frac{p_i^2}{m_i}, M = \sum_i m_i, q_1$

$$I = \sum_i m_i q_i^2 = M^{-1} \sum_{i < j} m_i m_j (q_i - q_j)^2$$

$$\therefore \frac{m_i m_j (q_i - q_j)^2}{MI} = \frac{m_k v_k^2}{K}$$

$$p_2$$

Oriented area (L=0, dI/dt=0)

$$p_1 \wedge p_2 = -k^2 (q_2 - q_1) \wedge (q_3 - q_1)$$

= $-\frac{K}{MI} m_1 m_2 m_3 (q_1 \wedge q_2 + q_2 \wedge q_3 + q_3 \wedge q_1)$
= $-\frac{K}{I} m_1 m_2 q_1 \wedge q_2.$

 $\therefore \sum_i m_i q_i = 0 \Rightarrow m_1 m_2 q_1 \land q_2 = m_2 m_3 q_2 \land q_3 = m_3 m_1 q_3 \land q_1.$

$$\frac{q_i \wedge q_j}{I} + \frac{v_i \wedge v_j}{K} = 0.$$

Energy balance for L=0, dI/dt=0 orbit under 1/r^2

$$\frac{d^2 I}{dt^2} = 0 \Rightarrow K = \sum_{i < j} \frac{m_i m_j}{r_{ij}^2}$$
$$L = 0, \quad \frac{dI}{dt} = 0 \Rightarrow \frac{1}{r_{ij}^2} = \frac{m_1 m_2 m_3 K}{MI} \quad \frac{1}{p_k^2}.$$

$$\therefore K = \frac{m_1 m_2 m_3 K}{MI} \left(\frac{m_1 m_2}{p_3^2} + \frac{m_2 m_3}{p_1^2} + \frac{m_3 m_1}{p_2^2} \right)$$

$$\therefore \frac{m_1 m_2}{p_3^2} + \frac{m_2 m_3}{p_1^2} + \frac{m_3 m_1}{p_2^2} = \frac{MI}{m_1 m_2 m_3} = \text{const.}$$

Geometrical properties of L=0 orbit

• In the followings, we consider L=0 orbit, not assuming dI/dt=0.

• Even in this case, we can find the synchronised similar triangles.

Synchronised similar triangles for L=0 orbit

For $L = \sum_{i} m_i q_i \wedge v_i = 0$ but $\frac{dI}{dt} = \sum_{i} m_i q_i \cdot v_i \neq 0$ orbit, consider

$$\xi_i = \frac{q_i}{\sqrt{I}}, \ \eta_i = \frac{d\xi_i}{dt} = \frac{v_i}{\sqrt{I}} - \frac{1}{2I}\frac{dI}{dt}\frac{q_i}{\sqrt{I}}, \ \mu_i = m_i.$$

Then, we have

$$\sum_{i} \mu_i \xi_i = 0, \sum_{i} \mu_i \eta_i = 0, \sum_{i} \mu_i \xi_i \wedge \eta_i = 0, \sum_{i} \mu_i \xi_i \cdot \eta_i = 0.$$

 \therefore Triangle whose vertexes are $\xi_i = \frac{q_i}{\sqrt{I}}$ and triangle whose perimeters are $\mu_i \eta_i = \mu_i \frac{d\xi_i}{dt}$ are always inversely similar. (Synchronised Similar Triangles)

Purely algebraic derivation of synchronised similar triangles

Let $i = 1, 2, 3, \xi_i, \eta_i \in \mathbb{R}^2, \mu_i > 0$ such that

$$\sum_{i} \mu_i \xi_i = 0, \sum_{i} \mu_i \eta_i = 0, \sum_{i} \mu_i \xi_i \wedge \eta_i = 0, \sum_{i} \mu_i \xi_i \cdot \eta_i = 0.$$

Let
$$M = \sum_{i} \mu_{i}$$
, $I(\xi) = \sum_{i} \mu_{i} \xi_{i}^{2} = M^{-1} \sum_{i < j} \mu_{i} \mu_{j} (\xi_{i} - \xi_{j})^{2}$. Then

$$\frac{\mu_k \xi_k^2}{I(\xi)} = \frac{\mu_i \mu_j (\eta_i - \eta_j)^2}{MI(\eta)}, \quad \frac{\mu_k \eta_k^2}{I(\eta)} = \frac{\mu_i \mu_j (\xi_i - \xi_j)^2}{MI(\xi)},$$
$$\frac{\mu_k \xi_k^2}{I(\xi)} + \frac{\mu_k \eta_k^2}{I(\eta)} = \frac{\mu_i \mu_j (\xi_i - \xi_j)^2}{MI(\xi)} + \frac{\mu_i \mu_j (\eta_i - \eta_j)^2}{MI(\eta)} = \frac{\mu_i + \mu_j}{M}$$

and

$$\frac{\xi_i \wedge \xi_j}{I(\xi)} + \frac{\eta_i \wedge \eta_j}{I(\eta)} = 0.$$

Synchronised similar triangles for L=0 orbit

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Therefore, we get

$$\frac{\xi_i \wedge \xi_j}{I(\xi)} + \frac{\eta_i \wedge \eta_j}{I(\eta)} = 0$$

where

$$\xi_i = \frac{q_i}{\sqrt{I}}, \ \eta_i = \frac{d\xi_i}{dt} = \frac{v_i}{\sqrt{I}} - \frac{1}{2I}\frac{dI}{dt}\frac{q_i}{\sqrt{I}}$$

That is

$$\frac{q_i \wedge q_j}{I} + \frac{v_i \wedge v_j}{K} = \frac{1}{2IK} \frac{dI}{dt} \frac{d}{dt} (q_i \wedge q_j)$$

where

$$I = \sum_{i} m_i q_i^2, \ K = \sum_{i} m_i v_i^2.$$

Evolution of oriented area for L=0 orbit

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Let

 $\Delta = \frac{1}{2}(q_2 - q_1) \wedge (q_3 - q_1) = \sum_{i < j} (q_i \wedge q_j).$

Then

$$\frac{d^2}{dt^2}(q_i \wedge q_j) = \frac{d^2 q_i}{dt^2} \wedge q_j + q_i \wedge \frac{d^2 q_j}{dt^2} + 2v_i \wedge v_j.$$

$$I = -\frac{K}{I}(q_i \wedge q_j) + \frac{1}{2I}\frac{dI}{dt}\frac{d}{dt}(q_i \wedge q_j)$$

$$\therefore I \frac{d}{dt}\left(\frac{1}{I}\frac{d\Delta}{dt}\right) = -\left(\frac{2K}{I} + \sum_{i < j}(m_i + m_j)r_{ij}^{\alpha - 2}\right)\Delta.$$

Infinitely many syzygies or collisions (Montgomery 2002)

- Montgomery formulated and proved: *Any bounded three-body orbit with L=0 has infinitely many collinear configurations (syzygies or eclipses) or collisions.*
- We give a simple proof.

A simple proof of infinitely many syzygies or collisions

The same in the ball to an a state of the a flat wing

Proof. Let

$$S(t) = \frac{\Delta(t)}{\sqrt{I(t)}}.$$

Then

$$\frac{d^2 S}{dt^2} = -\left\{ \sum_{i < j} (m_i + m_j) r_{ij}^{\alpha - 2} + \frac{2K}{I} + \frac{1}{2I} \frac{d^2 I}{dt^2} - \frac{3}{4I^2} \left(\frac{dI}{dt}\right)^2 \right\} S$$
$$= -\omega^2 S,$$

$$\omega^2 \ge \omega_0^2 > 0$$
 where $\omega_0^2 = M\left(\frac{m_{min}^2}{MI_{max}}\right)^{(2-\alpha)/2}$ for $\alpha \le 2$.

Conclusion 1: Synchronised similar triangles for L=0, dI/dt=0 orbit.



Conclusion 2: Synchronised similar triangles for figure-eight under 1/r^2



$$q'_i = \frac{q_i}{\sqrt{I}} \qquad \qquad m_i = 1 \qquad \qquad p'_i = \frac{p_j - p_k}{\sqrt{3K}}$$

Two triangles are inversely congruent.

$$\sum_{i} \frac{1}{p_i^2} = 3I$$

Conclusion 3: Synchronised similar triangles for L=0 orbit.

The variables form the synchronised similar triangles

$$\xi_i = \frac{q_i}{\sqrt{I}}, \ \eta_i = \frac{d\xi_i}{dt} = \frac{v_i}{\sqrt{I}} - \frac{1}{2I}\frac{dI}{dt}\frac{q_i}{\sqrt{I}}.$$

Then, we get

$$\frac{q_i \wedge q_j}{I} + \frac{v_i \wedge v_j}{K} = \frac{1}{2IK} \frac{dI}{dt} \frac{d}{dt} (q_i \wedge q_j),$$
$$I\frac{d}{dt} \left(\frac{1}{I}\frac{d\Delta}{dt}\right) = -\left(\frac{2K}{I} + \sum_{i < j} (m_i + m_j)r_{ij}^{\alpha - 2}\right) \Delta.$$

We gave a short proof that $\Delta(t)$ has infinitely many zeros if $\alpha \leq 2$.

I have a dream. One day, someone will e-mail me

and say

"I have *solved* the figure-eight !"

Thank you.